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GENERALIZED K_4 -FUNCTION AND ITS APPLICATION IN SOLVING KINETIC EQUATION OF FRACTIONAL ORDER

AHMAD FARAJ¹, TARIQ SALIM^{1,*}, SAFAA SADEK², JAMAL ISMAIL²

¹Department of Mathematics, Al-Azhar University – Gaza. P.O. Box 1277. Gaza – Palestine

²Department of mathematics, College of Girls Ain Shams University – Cairo – Egypt

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Abstract. This paper is devoted to introduce a new generalized K_4 - function in terms of some special functions. The differ-integration of this function is also investigated. A method for deriving the solution of the generalized fractional kinetic equation in term of the generalized K_4 - function defined, generalized M-Series $M_{p,q,m,n}^{\alpha,\beta}(z)$ and generalized Mittag – Leffler function $E_{\alpha,\beta,q}^{\gamma,\delta,p}(z)$ is investigated. The applied method depends on the fractional differ-integral operator techniques.

Keywords: generalized K_4 -function; generalized M-series; generalized Mittag Leffler function; fractional kinetic equation; fractional differintegral operator.

2010 AMS Subject Classification: 82B40.

1. Introduction

Fractional calculus is a field that deals with derivative and integral of arbitrary orders which almost used at every field of mathematics namely special functions. The Mittag – Leffler function has gained importance during the last century due to its applications in the solution of fractional order differential and integral equations, that function is introduced by Mittag – Leffler [8] in terms of power series

*Corresponding author

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$$E_\alpha(z) = \sum_{k=0}^{\infty} \frac{z^k}{\Gamma(\alpha k + 1)} \quad \alpha > 0 \tag{1.1}$$

Many authors defined and studied in their research papers different generalization of Mittag – Leffler type function like $E_{\alpha,\beta}(z)$ defined by Wiman[25], $E_{\alpha,\beta}^\gamma(z)$ studied by Prabahaker [11], $E_{\alpha,\beta}^{\gamma,q}(z)$ introduced by Shukla and Prajapati [24] and $E_{\alpha,\beta}^{\gamma,\delta}(z)$ investigated by Salim [14].

A new generalization of Mittag – Leffler type function introduced by Salim and Faraj [15] as

$$E_{\alpha,\beta,q}^{\gamma,\delta,p}(z) = \sum_{k=0}^{\infty} \frac{(\gamma)_{pk} z^k}{\Gamma(\alpha k + \beta)(\delta)_{qk}} \tag{1.2}$$

where $z, \alpha, \beta, \gamma, \delta \in \mathbb{C}$, $\min\{\text{Re}(\alpha), \text{Re}(\beta), \text{Re}(\gamma), \text{Re}(\delta)\} > 0$, $p, q > 0$

The authors [1] introduced in a recent paper a new generalization of M-Series $M_{p,q,m,n}^{\alpha,\beta}(z)$ as

$$M_{p,q,m,n}^{\alpha,\beta}(z) = \sum_{k=0}^{\infty} \frac{(a_1)_{km}, \dots, (a_p)_{km}}{(b_1)_{kn}, \dots, (b_q)_{kn}} \frac{z^k}{\Gamma(\alpha k + \beta)} \tag{1.3}$$

where $z, \alpha, \beta \in \mathbb{C}$, m, n nonnegative real number, also of the parameter b_j, s is negative or zero.

The Series in (1.3) is defined depending on the M-Series $M_{p,q}^\alpha(z)$ introduced by Sharma [22] and its generalized M-Series $M_{p,q}^{\alpha,\beta}(z)$ studied by Sharma and Jain [23].

The new generalization of the M-Series (1.3) is interesting because the ${}_pF_q(z)$ hypergeometric function and generalized Mittag – Leffler function (1.2) follow as its particular cases [2,15].

The interest R and G -function defined by Lorenzo and hartely [5], [6], and their populanty have sharply increased in view of their importance role and applications in fractional calculus.

$$R_{\alpha,\beta}[a,c,z] = \sum_{k=0}^{\infty} \frac{(a)^k (z-c)^{(k+1)\alpha-\beta-1}}{\Gamma(\alpha(k+1)-\beta)} \tag{1.4}$$

and

$$G_{\alpha,\beta,\gamma}[a,c,z] = \sum_{k=0}^{\infty} \frac{(\gamma)_k (a)^k (z-c)^{(k+\gamma)\alpha-\beta-1}}{k! \Gamma((k+\gamma)\alpha-\beta)} \tag{1.5}$$

Recently Sharma [20] defined K_4 -function

$$K_4^{(\alpha, \beta, \gamma); (a, c); (p, q)}(z) = \sum_{k=0}^{\infty} \frac{(a_1)_k, \dots, (a_p)_k}{(b_1)_k, \dots, (b_q)_k} \frac{(\gamma)_k (a)^k (z-c)^{(k+\gamma)\alpha-\beta-1}}{\Gamma((k+\gamma)\alpha-\beta)} \quad (1.6)$$

which is closely related to another special functions especially the R and G -function and M-Series defined by Sharma and Jain [23].

On the other hand, Fractional kinetic equation have gained importance due to their occurrence in science and engineering, the generalized fractional kinetic equation in term of Mittag – Leffler function studied by Sexcena, Mathai and Haubold [19], they introduced the solution of the generalized fractional kinetic equation associated with generalized Mittag – Leffler function and the R -function, for more result one can refer to the work of Sharma [21], Saichev and Zaslavsky [13], Sexcena [18], Zaslavsky [26], and Sexcena, Kalla [17].

Recently, Gupta and Parihar [3], introduced an alternative method for solving generalized fractional kinetic equation involving the generalized functions for the fractional calculus based on fractional differ-integral operator technique which differ from Laplace transform operator method.

This paper is divided to:

- Define a new generalized K_4 -function and its relation to the other special functions.
- Investigate the differ-integration propriety of the new function.
- Solving Fractional and general fractional kinetic equation in terms of the new generalized K_4 -function, the generalized M-Series and the generalized Mittag – Leffler function.

2. A new special function

The generalized K_4 -function introduced by the authors is defined as follows.

$$K_{4(m, n)}^{(\alpha, \beta, \gamma); (a, c); (p, q)}(a_1, \dots, a_p, b_1, \dots, b_n; z) = K_{4(m, n)}^{(\alpha, \beta, \gamma); (a, c); (p, q)}(z) \\ = \sum_{k=0}^{\infty} \frac{(a_1)_{km}, \dots, (a_p)_{km}}{(b_1)_{kn}, \dots, (b_n)_{kn}} \frac{(\gamma)_k (a)^k (z-c)^{(k+\gamma)\alpha-\beta-1}}{k! \Gamma((k+\gamma)\alpha-\beta)} \quad (2.1)$$

where $\text{Re}(\alpha\gamma - \beta) > 0$ and $(a_i)_k$ ($i = 1, 2, \dots, p$) and $(b_j)_k$ ($j = 1, 2, \dots, q$) are the Pochhammer symbols.

The series (2.1) is defined when non of the parameters b_j 's is a negative integer or zero. If any numerator parameter a_i is anegative integer or zero, then the series terminate to a polynomial of z .

From the ratio test it is evident that the series is convergent for all z if $pm < qn + \text{Re}(\alpha)$, also when

$pm = qn + \text{Re}(\alpha)$ it is convergent in some cases, let $\xi = \sum_{j=1}^{pm} a_j - \sum_{j=1}^{qn} b_j$. It can be shown that when

$pm = qn + \text{Re}(\alpha)$, the series is absolutely convergent for $|z|=1$ if $\text{Re}(\xi) < 0$, conditionally convergent for $z = -1$ if $0 \leq \text{Re}(\xi) < 1$ and divergent for $|x|=1$ if $\text{Re}(\xi) \geq 1$.

Relation with another special functions:

- (i) Putting $m = n = 1$ in the generalized K_4 -function (2.1) becomes the K_4 -function (1.6) defined by Sharma [21]

$$K_{4(1,1)}^{(\alpha,\beta,\gamma),(a,c),(p,q)}(z) = K_4^{(\alpha,\beta,\gamma),(a,c),(p,q)}(z) \tag{2.2}$$

- (ii) When there is no upper and lower parameters of (2.1), we get

$$K_{4(m,n)}^{(\alpha,\beta,\gamma),(a,c),(0,0)}(-, -; z) = \sum_{k=0}^{\infty} \frac{(\gamma)_k (a)^k (z - c)^{(k+\gamma)\alpha - \beta - 1}}{k! \Gamma((k + \gamma)\alpha - \beta)} \tag{2.3}$$

which reduces to the G -function (1.5) defined by Lorenzo and Hartley [5] devoted by $G_{\alpha,\beta,\gamma}[a,c,z]$.

- (iii) If we put $\gamma = 1$ in (2.3).

$$K_{4(m,n)}^{(\alpha,\beta,1),(a,c),(0,0)}(-, -; z) = \sum_{k=0}^{\infty} \frac{(a)^k (z - c)^{(k+1)\alpha - \beta - 1}}{\Gamma((k + 1)\alpha - \beta)} \tag{2.4}$$

which reduces to the R -function (1.4) defined by Lorenzo and Hartley [6] and denoted by $R_{\alpha,\beta}[a,c,z]$.

- (iv) If we take $c = \beta = 0$ in (2.4).

$$K_{4(m,n)}^{(\alpha,0,1),(a,0),(0,0)}(-, -; z) = \sum_{k=0}^{\infty} \frac{a^k z^{(k+1)\alpha - 1}}{\Gamma((k + 1)\alpha)} \tag{2.5}$$

Which reduces to the F -function defined by Lorenzo and Hartley [5] and denoted by $F_{\alpha}[a,z]$.

Relation between generalized K_4 -function and generalized M-Series.

Putting $\beta = \alpha - \beta$, $\gamma = 1$, $a = 1$ and $c = 0$ in (2.1) we obtain

$$\begin{aligned} K_{4(m,n)}^{(\alpha, \alpha-\beta, 1), (1, 0), (p, q)}(a_1, \dots, a_p, b_1, \dots, b_q; z) &= \sum_{k=0}^{\infty} \frac{(a_1)_{km} \dots (a_p)_{km}}{(b_1)_{kn} \dots (b_q)_{kn}} \frac{(1)_k (1)_k (z)^{(k+1)\alpha - (\alpha-\beta) - 1}}{k! \Gamma((k+1)\alpha - (\alpha-\beta))} \\ &= z^{\beta-1} \sum_{k=0}^{\infty} \frac{(a_1)_{km} \dots (a_p)_{km}}{(b_1)_{kn} \dots (b_q)_{kn}} \frac{z^{k\alpha}}{\Gamma(k\alpha + \beta)} \\ &= z^{\beta-1} M_{p,q,m,n}^{\alpha,\beta}(a_1, \dots, a_p, b_1, \dots, b_q, z^\alpha) \end{aligned} \quad (2.6)$$

also if we put $p = q = 1$ in (2.6) we get

$$K_{4(m,n)}^{(\alpha, \alpha-\beta, 1), (1, 0), (1, 1)}(a_1, b_1; z) = z^{\beta-1} \sum_{k=0}^{\infty} \frac{(a_1)_{km}}{(b_1)_{kn}} \frac{(z^\alpha)^k}{\Gamma(\alpha k + \beta)} = z^{\beta-1} E_{m,n,1}^{\alpha,\beta,1}(z^\alpha)$$

3. Differ-integration of generalized K_4 -function

In this section we will derive the relation between generalized K_4 -function and the operator of differ-integral given by Oldham and Spainer [9]. The relation is presented in the next theorem as follows:

Theorem 3.1:

Let $-\infty < r < \infty$, $\operatorname{Re}(\alpha\gamma - \beta) > 0$, $z > c \geq 0$ and ${}_c d_z^r$ be the operator of differintegral given by Oldham and Spainer then the relation holds:

$${}_c d_z^r K_{4(m,n)}^{(\alpha, \beta, \gamma), (a, c), (p, q)}(a_1, \dots, a_p, b_1, \dots, b_q; z) = K_{4(m,n)}^{(\alpha, \beta+r, \gamma), (a, c), (p, q)}(a_1, a_2, \dots, a_p, b_1, \dots, b_q; z) \quad (3.1)$$

Proof:

The differ-integral operator defined by Oldham and Springer [9] of function $f(z)$ is given by

$${}_c d_z^r f(z) = \frac{d f(z)}{d(z-c)^r} \quad (3.2)$$

Using (2.1) and (3.2) we have

$${}_c d_z^r \{K_{4(m,n)}^{(\alpha, \beta, \gamma), (a, c), (p, q)}(z)\} = {}_c d_z^r \left\{ \sum_{k=0}^{\infty} \frac{(a_1)_{km} \dots (a_p)_{km}}{(b_1)_{kn} \dots (b_q)_{kn}} \frac{(\gamma)_k (a)^k}{k!} \frac{(z-c)^{(k+\gamma)\alpha - \beta - 1}}{\Gamma((k+\gamma)\alpha - \beta)} \right\}$$

$$\begin{aligned}
 &= \sum_{k=0}^{\infty} \frac{(a_1)_{km} \dots (a_p)_{km}}{(b_1)_{kn} \dots (b_q)_{kn}} \frac{(\gamma)_k (a)^k}{k! \Gamma((k+\gamma)\alpha-\beta)} {}_c d_z^r (z-c)^{(k+\gamma)\alpha-\beta-1} \\
 &= \sum_{k=0}^{\infty} \frac{(a_1)_{km} \dots (a_p)_{km}}{(b_1)_{kn} \dots (b_q)_{kn}} \frac{(\gamma)_k (a)^k}{k! \Gamma((k+\gamma)\alpha-\beta)} \\
 &\cdot ((k+\gamma)\alpha-\beta-1)((k+\gamma)\alpha-\beta-2)\dots((k+\gamma)\alpha-\beta-r)(z-c)^{(k+\gamma)\alpha-\beta-r-1} \\
 &= \sum_{k=0}^{\infty} \frac{(a_1)_{km} \dots (a_p)_{km}}{(b_1)_{kn} \dots (b_q)_{kn}} \frac{(\gamma)_k (a)^k}{k! \Gamma((\alpha+k)\gamma-(\beta+r))} (z-c)^{(k+\gamma)\alpha-(\beta+r)-1} \\
 &= K_{4(m,n)}^{(\alpha,\beta+r,\gamma),(a,c),(p,q)}(a_1, a_2, \dots, a_p, b_1, \dots, b_q; z)
 \end{aligned}$$

This shows that the differintegral of the generalized K_4 -function is again a generalized K_4 -function with indices $\beta+r$.

Particular case:

(1) ${}_c d_z^r G_{\alpha,\beta,\gamma}[a,c,z] = G_{\alpha,\beta+r,\gamma}[a,c,z]$

(2) ${}_c d_z^r R_{\alpha,\beta}[a,c,z] = R_{\alpha,\beta+r}[a,c,z]$

(3) ${}_0 d_z^r F_{\alpha}[a,0,z] = R_{\alpha+r}[a,0,z]$

4. Solving general fractional kinetic equation in terms of generalized K_4 -function, generalized M-Series and generalized Mittag – Leffler function

The generalized Riemann – Liouville operators of fractional calculus [7,9] are defined as

$${}_a D_z^{-\nu} f(z) = \frac{1}{\Gamma(\nu)} \int_a^z (z-u)^{\nu-1} f(u) du \quad \text{Re}(\nu) > 0, \quad z > a \tag{4.1}$$

with

${}_a D_z^0 f(z) = f(z)$, and

$${}_a D_z^{\mu} f(z) = \frac{d^k}{dz^k} ({}_a D_z^{\mu-k} f(z)) \quad \text{Re}(\mu) > 0, \quad k - \mu > 0 \tag{4.2}$$

If $p(z) = (z - a)^p$, we have from [10]

$${}_a D_z^{-\nu} (z - a)^{p-1} = \frac{\Gamma(p)}{\Gamma(p + \nu)} (z - a)^{p-\nu-1} \quad (4.3)$$

where $\operatorname{Re}(\nu) > 0$, $\operatorname{Re}(p) > 0$, $z > a$.

On integrating the standard kinetic equation

$$\frac{d}{dt} N_i(t) = -c_i N_i(t) \quad c_i > 0 \quad (4.4)$$

Haubold and Mathai [4] derived that

$$N_i(t) - N_i(0) = -c_i {}_0 D_t^{-1} N_i(t) \quad (4.5)$$

Where ${}_0 D_t^{-1}$ is the standard Riemann integral operator, $N_i = N_i(t)$ is number density of species i , which is a function of time t and $N_i(0) = N_0$ is the number density of that species of time $t = 0$.

By dropping the index i and replacing the Riemann integral ${}_0 D_t^{-1}$ operator by the fractional Riemann – Liouville operator ${}_0 D_t^{-\nu}$, the kinetic equation (4.5) is reduced to

$$N(t) - N_0 = -c^\nu {}_0 D_t^{-\nu} N(t) \quad (4.6)$$

Multiplying both side of (4.6) by $(-c^\nu)^r {}_0 D_t^{-r\nu}$ and taking the sum over r from 0 to ∞ yields

$$\sum_{r=0}^{\infty} (-c^\nu)^r {}_0 D_t^{-r\nu} N(t) - \sum_{r=0}^{\infty} (-c^\nu)^{r+1} {}_0 D_t^{-(r+1)\nu} N(t) = N_0 \sum_{r=0}^{\infty} (-c^\nu)^r {}_0 D_t^{-r\nu} 1$$

Replacing r by $(r-1)$ in the second sum of above equation and then cancelling the equal terms on the left hand side, and applying (4.3) after putting $p=1$ on the right hand side of above equation

$$N(t) = N_0 \sum_{r=0}^{\infty} (-c^\nu)^r {}_0 D_t^{-r\nu} 1 \quad (1)$$

Or

$$N(t) = N_0 \sum_{r=0}^{\infty} \frac{(-c^\nu)^r (t)^{-r\nu}}{\Gamma(1+r\nu)} = N_0 \sum_{r=0}^{\infty} \frac{(-c^\nu t^\nu)^r}{\Gamma(1+r\nu)}$$

Hence,

$$N(t) = N_0 M_{0,0}^{\nu,1}(-, -; -c^\nu t^\nu)$$

Theorem 4.1:

If $a, \beta, \alpha, > 0 ; \nu > 0$, then the solution of the general fractional kinetic equation

$$N(t) - N_0 t^{\beta-1} E_{\alpha, \beta, n}^{a, b, m}(at^\alpha) = -c^\nu {}_0D_t^{-\nu} N(t) \tag{4.7}$$

is given by

$$N(t) = N_0 \sum_{r=0}^{\infty} (-c^\nu)^r K_{4(m, n)}^{(\alpha, \alpha-\beta-r\nu, 1), (a, 0), (1, 1)}(t) \tag{4.8}$$

Proof:

Multiplying both side of (4.7) by $(-c^\nu)^r {}_0D_t^{-r\nu}$ and taking the sum over r from 0 to ∞ , we get

$$\begin{aligned} & \sum_{r=0}^{\infty} (-c^\nu)^r {}_0D_t^{-r\nu} N(t) - \sum_{r=0}^{\infty} (-c^\nu)^{r+1} {}_0D_t^{-(r+1)\nu} N(t) \\ &= N_0 \sum_{r=0}^{\infty} (-c^\nu)^r {}_0D_t^{-r\nu} t^{\beta-1} E_{\alpha, \beta, n}^{a, b, m}(at^\alpha) \end{aligned}$$

Replacing r by $(r-1)$ in the second sum of above equation and then cancelling the equal terms on the left hand side,

$$N(t) = N_0 \sum_{r=0}^{\infty} (-c^\nu)^r {}_0D_t^{-r\nu} t^{\beta-1} E_{\alpha, \beta, n}^{a, b, m}(at^\alpha)$$

Now

$${}_0D_t^{-r\nu} t^{\beta-1} E_{\alpha, \beta, n}^{a, b, m}(at^\alpha) = \sum_{k=0}^{\infty} \frac{(a)_{km}}{(b)_{kn}} \frac{(a)^k}{\Gamma(\alpha k + \beta)} {}_0D_t^{-r\nu} t^{\alpha k + \beta - 1}$$

and by applying (4.3), we get

$$\begin{aligned} & {}_0D_t^{-r\nu} t^{\beta-1} E_{\alpha, \beta, n}^{a, b, m}(at^\alpha) = \sum_{k=0}^{\infty} \frac{(a)_{km}}{(b)_{kn}} \frac{(a)^k}{\Gamma(\alpha k + \beta)} \frac{\Gamma(\alpha k + \beta)}{\Gamma(\alpha k + \beta + r\nu)} t^{r\nu + \alpha k + \beta - 1} \\ &= \sum_{k=0}^{\infty} \frac{(a)_{km}}{(b)_{kn}} \frac{(a)^k}{k!} \frac{(1)_k}{\Gamma(\alpha(k+1) - (\alpha - \beta - r\nu))} t^{(k+1)\alpha - (\alpha - \beta - r\nu) - 1} \end{aligned}$$

Then the solution become

$$N(t) = N_0 \sum_{r=0}^{\infty} (-c^\nu)^r K_{4(m, n)}^{(\alpha, \alpha-\beta-r\nu, 1), (a, 0), (1, 1)}(t)$$

Theorem4.2:

If $\alpha, \beta, a > 0$ and $\nu > 0$, then the solution of the general fractional kinetic equation

$$N(t) - N_0 t^{\beta-1} M_{p,q,m,n}^{\alpha,\beta}(at^\alpha) = -c^\nu {}_0D_t^{-\nu} N(t) \quad (4.9)$$

is given by

$$N(t) = N_0 \sum_{r=0}^{\infty} (-c^\nu)^r K_{4(m,n)}^{(\alpha,\alpha-\beta-r\nu,1),(a,0),(p,q)}(t) \quad (4.10)$$

Proof:

Multiplying both side of (4.9) by $(-c^\nu)^r {}_0D_t^{-r\nu}$ and taking the sum over r from 0 to ∞ , we get

$$\begin{aligned} & \sum_{r=0}^{\infty} (-c^\nu)^r {}_0D_t^{-r\nu} N(t) - \sum_{r=0}^{\infty} (-c^\nu)^{r+1} {}_0D_t^{-(r+1)\nu} N(t) \\ &= N_0 \sum_{r=0}^{\infty} (-c^\nu)^r {}_0D_t^{-r\nu} t^{\beta-1} M_{p,q,m,n}^{\alpha,\beta}(at^\alpha) \end{aligned}$$

Using the same technique of theorem (4.1), $N(t)$ become

$$N(t) = N_0 \sum_{r=0}^{\infty} (-c^\nu)^r {}_0D_t^{-r\nu} t^{\beta-1} M_{p,q,m,n}^{\alpha,\beta}(at^\alpha)$$

Now

$$\begin{aligned} {}_0D_t^{-r\nu} t^{\beta-1} M_{p,q,m,n}^{\alpha,\beta}(at^\alpha) &= \sum_{k=0}^{\infty} \frac{(a_1)_{km} \dots (a_p)_{km} (a)^k {}_0D_t^{-r\nu} (t^{\alpha k + \beta - 1})}{(b_1)_{kn} \dots (b_q)_{kn} \Gamma(\alpha k + \beta)} \\ &= \sum_{k=0}^{\infty} \frac{(a_1)_{km} \dots (a_p)_{km} (a)^k}{(b_1)_{kn} \dots (b_q)_{kn}} \frac{\Gamma(\alpha k + \beta)}{\Gamma(\alpha k + \beta + r\nu)} t^{r\nu + \alpha k + \beta - 1} \\ &= \sum_{k=0}^{\infty} \frac{(a_1)_{km} \dots (a_p)_{km} (1)_k (a)^k}{(b_1)_{kn} \dots (b_q)_{kn} k!} \frac{t^{(k+1)\alpha - (\alpha - \beta - r\nu) - 1}}{\Gamma((k+1)\alpha - (\alpha - \beta - r\nu))} \\ &= K_{4(m,n)}^{(\alpha,\alpha-\beta-r\nu,1),(a,0),(p,q)}(t) \end{aligned}$$

So that the solution become

$$N(t) = N_0 \sum_{r=0}^{\infty} (-c^\nu)^r K_{4(m,n)}^{(\alpha,\alpha-\beta-r\nu,1),(a,0),(p,q)}(t).$$

Theorem 4.3:

If $a, c, \gamma, \alpha, \beta, a > 0$; $\nu > 0$ and $\text{Re}(\gamma\alpha - \beta) > 0$, then the equation

$$N(t) - N_0 K_{4(m,n)}^{(\alpha,\beta,\gamma),(a,b),(p,q)}(t) = -c^\nu {}_0D_t^{-\nu} N(t) \quad (4.11)$$

has the solution

$$N(t) = N_0 \sum_{r=0}^{\infty} (-c^v)^r K_{4(m,n)}^{(\alpha, \beta-rv, \gamma); (a,b), (p,q)}(t) \tag{4.12}$$

Proof:

Repeating the process applied in the theorem (4.1), (4.2), we get

$$\begin{aligned} & \sum_{r=0}^{\infty} (-c^v)^r {}_0D_t^{-rv} N(t) - \sum_{r=0}^{\infty} (-c^v)^{r+1} {}_0D_t^{-(r+1)v} N(t) \\ &= N_0 \sum_{r=0}^{\infty} (-c^v)^r {}_0D_t^{-rv} K_{4(m,n)}^{(\alpha, \beta, \gamma), (a,c), (p,q)}(t) \end{aligned}$$

Hence $N(t)$ become

$$N(t) = N_0 \sum_{r=0}^{\infty} \sum_{k=0}^{\infty} \frac{(-c^v)^r (a_1)_{km} \dots (a_p)_{km} (\gamma)_k (a)^k {}_0D_t^{-rv} (t-b)^{(k+\gamma)\alpha-\beta-1}}{(b_1)_{kn} \dots (b_q)_{kn} k! \Gamma((k+\gamma)\alpha-\beta)}$$

and by applying (4.3), the solution become

$$N(t) = N_0 \sum_{r=0}^{\infty} \sum_{k=0}^{\infty} \frac{(-c^v)^r (a_1)_{km} \dots (a_p)_{km} (\gamma)_k (a)^k \Gamma((k+\gamma)\alpha-\beta) (t-b)^{(k+\gamma)\alpha-\beta+rv-1}}{(b_1)_{kn} \dots (b_q)_{kn} k! \Gamma((k+\gamma)\alpha-\beta) \Gamma((k+\gamma)\alpha-\beta+rv)}$$

Or

$$N(t) = N_0 \sum_{r=0}^{\infty} (-c^v)^r K_{4(m,n)}^{(\alpha, \beta-rv, \gamma), (a,c), (p,q)}(t).$$

Conflict of Interests

The author declares that there is no conflict of interests.

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