



Available online at <http://scik.org>

J. Math. Comput. Sci. 2 (2012), No. 5, 1335-1352

ISSN: 1927-5307

ESTIMATION OF EIGEN FUNCTIONS TO THE NEW TYPE OF SPECTRAL PROBLEM

KARWAN H.F JWAMER*, ARYAN ALI. M

Department of Mathematics, Faculty of Science and Science Education ,School of Science, University of Sulaimani, Kurdistan Region, Sulaimani, Iraq

Abstract:In this paper, we study some properties of eigenvalues and the corresponding eigen functions of new type of spectral problem (1)-(4).

Keywords: Spectral problem, eigenvalues, eigen functions.

2010 AMS Subject Classification: 47E05, 34B05, 34B07

1. Introduction

In this paper, we study the new type of spectral problem T_0 which is defined by:

$$-y''(x) + y'(x) = \lambda^2 \rho(x)y(x), x \in [0, a], \quad (1)$$

$$y(0) = y'(0) = y(a) + y'(a) = 0, \quad (2)$$

$$\int_0^a y'(x)\bar{y}(x)dx = \tau^2 \quad (\tau \text{ is constant}), \quad (3)$$

$$\left(\int_0^a \rho(x)|y(x)|^2 dx \right)^{\frac{1}{2}} = 1, \quad (4)$$

where λ is a spectral parameter, and $\lambda = \delta + i\sigma$, where $\delta, \sigma \in \mathbb{R}$, and $i = \sqrt{-1}$. Let $a > 0$, we assume that $\rho(x) = \rho$ is a constant and let m and M be fixed such that $0 < m \leq M$. Let

* Corresponding author

Received May 6, 2012

$V^+[0, a]$ denotes the family of all positive integrable functions $\rho(x)$ on the closed interval $[0, a]$ that satisfy the condition $0 < m \leq \rho(x) \leq M$, equipped with usual L_1 metric. In what follows we refer to these functions as weight functions. Here we attempt to specify the properties of eigenvalues of the spectral problem T_o and estimating the eigen functions corresponding to the eigenvalues. For the first time an Italian physicist T. Regge [7] has studied the differential equation $-y'' + q(x)y = \lambda^2 \rho(x)y(x)$, $x \in (0, a)$ with the boundary condition $y(0) = 0, y'(a) - i\lambda y(a) = 0$, and was considered by who showed that the system of eigen functions of this problem are completed and studied asymptotic behavior of eigenvalues of this problem $\rho(x) = 1$. Kravitsky [6] specified a class of functions that allowed expansion in uniformly convergent series in eigen functions and associated functions in the Regge problem when $\rho(x) \equiv 1$. The present time they are many Arthurs studied the estimation of eigen functions to the equation $-y'' + q(x)y = \lambda^2 \rho(x)y(x)$ but with different boundary conditions for more known about their works see [1-5].

2. Features of Eigenvalues of the problem T_o

Here we determine the properties of the eigenvalues of our problem T_o with the given boundary conditions.

Theorem 1: *Let $y(x)$ be an eigen function corresponding to the eigenvalue λ of the problem T_o , and $\rho(x) = \rho$ is a constant, then: (i) If $\delta \neq 0$, then λ is real.*

(ii) If $\sigma \neq 0$, then λ is complex.

Proof: Multiplying equation (1) by $\bar{y}(x)$ and integrating the obtained equation

from 0 to a , yields:

$$-\int_0^a y''(x)\bar{y}(x)dx + \int_0^a y'(x)\bar{y}(x)dx = \lambda^2 \int_0^a \rho(x)y(x)\bar{y}(x)dx$$

$$-\bar{y}(x)y'(x) \Big|_0^a + \int_0^a y'(x)\bar{y}'(x)dx + \int_0^a y'(x)\bar{y}(x)dx = \lambda^2 \int_0^a \rho(x)|y(x)|^2 dx$$

$$-\bar{y}(a)y'(a) + \bar{y}(0)y'(0) + \int_0^a y'(x)\bar{y}'(x)dx + \int_0^a y'(x)\bar{y}(x)dx = \lambda^2 \int_0^a \rho(x)|y(x)|^2 dx$$

By using boundary conditions (2), we get:

$$\bar{y}(a)y(a) + \int_0^a |y'(x)|^2 dx + \int_0^a y'(x)\bar{y}(x)dx = \lambda^2 \int_0^a \rho(x)|y(x)|^2 dx$$

In view of condition (3) and normalized condition (4), we have:

$$|y(a)|^2 + \int_0^a |y'(x)|^2 dx + \tau^2 = \lambda^2 \quad (5)$$

From equation (1) and the conditions (2)-(4) replace $y(x)$ by $\bar{y}(x)$, we get:

$$-\bar{y}''(x) + \bar{y}'(x) = \bar{\lambda}^2 \rho(x)\bar{y}(x)$$

$$\bar{y}(0) = \bar{y}'(0) = \bar{y}(a) + \bar{y}'(a) = 0, \int_0^a \bar{y}'(x)y(x)dx = \tau^2.$$

Multiplying the above differential equation by $y(x)$ and integrate from 0 up to a , we obtain:

$$|y(a)|^2 + \int_0^a |y'(x)|^2 dx + \tau^2 = \bar{\lambda}^2 \quad (6)$$

Subtracting equation (6) from equation (5) yields:

$$\lambda^2 - \bar{\lambda}^2 = 0 \rightarrow (\lambda - \bar{\lambda})(\lambda + \bar{\lambda}) = 0, (\lambda - \bar{\lambda}) = 0 \text{ or } (\lambda + \bar{\lambda}) = 0, \text{ then:}$$

- (i) If $\delta \neq 0$, $\therefore (\lambda + \bar{\lambda}) \neq 0$, thus $(\lambda - \bar{\lambda}) = 0 \rightarrow \lambda = \bar{\lambda}$, then λ is real.
- (ii) If $\sigma \neq 0$, so $(\lambda - \bar{\lambda}) \neq 0$, hence $(\lambda + \bar{\lambda}) = 0 \rightarrow \lambda = -\bar{\lambda}$, then λ is complex.

3. Estimation of Eigen functions of problem T_o

In this section, we estimate the eigen function $y(x)$ corresponding to eigenvalue λ of problem T_o .

Theorem 2: Let λ be an eigenvalue corresponding to the eigen function $y(x)$ of problem T_o , and $\rho(x) \in L^+[0, a]$, and $\delta \neq 0$, then

$$\lim_{n \rightarrow \infty} \frac{\max_{x \in [0, a]} |y(x)|}{|\lambda|^{\frac{1}{2}}} = A, \text{ where } A = \frac{\sqrt{2}}{\sqrt[4]{m}}.$$

Proof:

Let us consider the identity:

$$\begin{aligned} |y(x)|^2 &= y(x)\bar{y}(x) = \int_0^x [\bar{y}(t)y'(t) + y(t)\bar{y}'(t)]dt + |y(0)|^2 \\ &= \int_0^x \frac{\sqrt{\rho(t)}[\bar{y}(t)y'(t) + y(t)\bar{y}'(t)]}{\sqrt{\rho(t)}} dt + |y(0)|^2 \end{aligned}$$

From inequality $(t) \geq m$, we get:

$$\begin{aligned} |y(x)|^2 &\leq \int_0^x \frac{\sqrt{\rho(t)}|\bar{y}(t)y'(t) + y(t)\bar{y}'(t)|}{\sqrt{m}} dt + |y(0)|^2 \\ &\leq \frac{1}{\sqrt{m}} \left[\int_0^x \sqrt{\rho(t)}|\bar{y}(t)y'(t)|dt + \int_0^x \sqrt{\rho(t)}|y(t)\bar{y}'(t)|dt \right] + |y(0)|^2 \\ &\leq \frac{1}{\sqrt{m}} \left[\int_0^x \sqrt{\rho(t)}|\bar{y}(t)||y'(t)|dt + \int_0^x \sqrt{\rho(t)}|y(t)||\bar{y}'(t)|dt \right] + |y(0)|^2 \\ &= \frac{2}{\sqrt{m}} \int_0^x \sqrt{\rho(t)}|y(t)||y'(t)| + |y(0)|^2 \end{aligned}$$

And from boundary condition (2), $y(0) = 0$, therefore

$$\begin{aligned} |y(x)|^2 &\leq \frac{2}{\sqrt{m}} \int_0^x \sqrt{\rho(t)} |y(t)| |y'(t)| \\ &\leq \frac{2}{\sqrt{m}} \int_0^a \sqrt{\rho(t)} |y(t)| |y'(t)|. \end{aligned}$$

Using Bunyakovsky's inequality on the last inequality, we shall obtain:

$$|y(x)|^2 \leq \frac{2}{\sqrt{m}} \left[\int_0^a \rho(t) |y(t)|^2 dt \right]^{\frac{1}{2}} \left[\int_0^a |y'(t)|^2 dt \right]^{\frac{1}{2}}$$

From normality condition (4) we have: $\left[\int_0^a \rho(t) |y(t)|^2 dt \right]^{\frac{1}{2}} = 1$, hence

$$|y(x)|^2 \leq \frac{2}{\sqrt{m}} \left[\int_0^a |y'(t)|^2 dt \right]^{\frac{1}{2}} \quad (7)$$

From equation (5), we have:

$$\int_0^a |y'(x)|^2 dx = \lambda^2 - |y(a)|^2 - \tau^2, \text{ therefore equation (7) becomes:}$$

$$|y(x)|^2 \leq \frac{2}{\sqrt{m}} [\lambda^2 - |y(a)|^2 - \tau^2]^{\frac{1}{2}} = \frac{2}{\sqrt{m}} [\lambda^2 - (|y(a)|^2 + \tau^2)]^{\frac{1}{2}}$$

And since $\delta \neq 0$, so by theorem (1) λ is real, hence $\lambda^2 = |\lambda|^2$, thus the last inequality becomes:

$$|y(x)|^2 \leq \frac{2}{\sqrt{m}} \left[|\lambda|^2 - (|y(a)|^2 + \tau^2) \right]^{\frac{1}{2}} = \frac{2|\lambda|}{\sqrt{m}} \left[1 - \frac{(|y(a)|^2 + \tau^2)}{|\lambda|^2} \right]^{\frac{1}{2}}$$

Or

$$|y(x)|^2 \leq \frac{2}{\sqrt{m}} |\lambda| \rightarrow |y(x)| \leq |\lambda|^{\frac{1}{2}} \sqrt{\frac{2}{\sqrt{m}}}$$

And since x is any value in the interval $[0, a]$, thus

$$\max_{x \in [0, a]} |y(x)| \leq |\lambda|^{\frac{1}{2}} \sqrt{\frac{2}{\sqrt{m}}} \rightarrow \frac{\max_{x \in [0, a]} |y(x)|}{|\lambda|^{\frac{1}{2}}} \leq \frac{\sqrt{2}}{\sqrt[4]{m}}$$

Hence

$$\lim_{n \rightarrow \infty} \frac{\max_{x \in [0, a]} |y(x)|}{|\lambda|^{\frac{1}{2}}} = A, \text{ where } A = \frac{\sqrt{2}}{\sqrt[4]{m}}.$$

Theorem 2.3.2: Let ρ be a constant in the problem T_0 and if $y(x)$ is an eigen function of the problem T_0 , then $y(x)$ satisfy the inequality

$$\frac{1}{\sqrt{|\lambda|}} K_1 \leq \max_{x \in [0, a]} |y(x)| \leq \frac{1}{\sqrt{|\lambda|}} K_2,$$

Where K_1 and K_2 are constants.

Proof:

From equation (1), we have $y''(x) - y'(x) + \lambda^2 \rho y(x) = 0$ this is second order linear differential equation with constant coefficients, and then general solution is:

$$y(x) = e^{\frac{1}{2}x} \left[c_1 e^{i\sqrt{\lambda^2 \rho - \frac{1}{4}}x} + c_2 e^{-i\sqrt{\lambda^2 \rho - \frac{1}{4}}x} \right]$$

Applying the condition $y(0) = 0$, yields $c_2 = -c_1$, then we have

$$y(x) = c_1 \left[e^{\left(\frac{1}{2} + i\sqrt{\lambda^2 \rho - \frac{1}{4}}\right)x} - e^{\left(\frac{1}{2} - i\sqrt{\lambda^2 \rho - \frac{1}{4}}\right)x} \right]$$

Then

$$y'(x) = c_1 \left[\left(\frac{1}{2} + i\sqrt{\lambda^2 \rho - \frac{1}{4}} \right) e^{\left(\frac{1}{2} + i\sqrt{\lambda^2 \rho - \frac{1}{4}}\right)x} - \left(\frac{1}{2} - i\sqrt{\lambda^2 \rho - \frac{1}{4}} \right) e^{\left(\frac{1}{2} - i\sqrt{\lambda^2 \rho - \frac{1}{4}}\right)x} \right]$$

From the boundary condition $y(a) + y'(a) = 0$, we obtain:

$$c_1 \left[e^{\left(\frac{1}{2}+i\sqrt{\lambda^2\rho-\frac{1}{4}}\right)a} - e^{\left(\frac{1}{2}-i\sqrt{\lambda^2\rho-\frac{1}{4}}\right)a} \right] +$$

$$c_1 \left[\left(\frac{1}{2} + i\sqrt{\lambda^2\rho - \frac{1}{4}}\right) e^{\left(\frac{1}{2}+i\sqrt{\lambda^2\rho-\frac{1}{4}}\right)a} - \left(\frac{1}{2} - i\sqrt{\lambda^2\rho - \frac{1}{4}}\right) e^{\left(\frac{1}{2}-i\sqrt{\lambda^2\rho-\frac{1}{4}}\right)a} \right] = 0$$

Dividing both sides of the above equation by c_1 , we get:

$$\frac{\frac{3}{2}-i\sqrt{\lambda^2\rho-\frac{1}{4}}}{\frac{3}{2}+i\sqrt{\lambda^2\rho-\frac{1}{4}}} = e^{2i\sqrt{\lambda^2\rho-\frac{1}{4}}a} \quad (8)$$

The resulting equation (8) is used for specifying the eigenvalues of our problem.

To find the coefficient c_1 , we use the normalization condition (4)

$$\int_0^a \rho |c_1|^2 \left| e^{\frac{1}{2}x} \left[e^{i\sqrt{\lambda^2\rho-\frac{1}{4}}x} - e^{-i\sqrt{\lambda^2\rho-\frac{1}{4}}x} \right] \right|^2 dx = 1,$$

Or

$$\rho |c_1|^2 \int_0^a \left| e^{\frac{1}{2}x} \left[e^{i\sqrt{\lambda^2\rho-\frac{1}{4}}x} - e^{-i\sqrt{\lambda^2\rho-\frac{1}{4}}x} \right] \right|^2 dx = 1.$$

We introduce the notation $\alpha + i\beta = i\sqrt{\lambda^2\rho - \frac{1}{4}}$ (where α and β are real numbers),

then

$$\rho |c_1|^2 \int_0^a \left| e^{\frac{1}{2}x} \left[e^{(\alpha+i\beta)x} - e^{-(\alpha+i\beta)x} \right] \right|^2 dx = 1,$$

Or

$$\rho |c_1|^2 \int_0^a \left| e^{\left(\frac{1}{2}+\alpha\right)x+i\beta x} - e^{\left(\frac{1}{2}-\alpha\right)x-i\beta x} \right|^2 dx = 1 \quad (9)$$

Since

$$\left| e^{\left(\frac{1}{2}+\alpha\right)x+i\beta x} - e^{\left(\frac{1}{2}-\alpha\right)x-i\beta x} \right|^2 = 2 e^x (\cosh 2\alpha x - \cos 2\beta x)$$

Thus equation (9) becomes:

$$\rho |c_1|^2 \int_0^a 2 e^x (\cosh 2\alpha x - \cos 2\beta x) dx = 1.$$

By integrating the last equation by parts, we obtain

$$2\rho |c_1|^2 \left[\frac{1}{2(1+2\alpha)} (e^{(1+2\alpha)a} - 1) + \frac{1}{2(1-2\alpha)} (e^{(1-2\alpha)a} - 1) \right. \\ \left. - \frac{(2\beta \sin 2\beta a + \cos 2\beta a)}{(4\beta^2 + 1)} e^a + \frac{2\beta}{(4\beta^2 + 1)} \right] = 1$$

After some algebraic operations, we get

$$|c_1|^2 = (1 - 4\alpha^2)(4\beta^2 + 1) / 2\rho [e^a (\cosh 2\alpha a - 2\alpha \sinh 2\alpha a)(4\beta^2 + 1) \\ - (4\beta^2 + 1) + 2\beta(1 - 4\alpha^2) - e^a (2\beta \sin 2\beta a + \cos 2\beta a)(1 - 4\alpha^2)]$$

Or

$$|c_1| = \frac{1}{\sqrt{2\rho}} \frac{1}{\sqrt{\frac{1}{(1-4\alpha^2)} [e^a (\cosh 2\alpha a - 2\alpha \sinh 2\alpha a) - 1] + \frac{1}{(4\beta^2+1)} [2\beta - e^a (2\beta \sin 2\beta a + \cos 2\beta a)]}}$$

By substituting $|c_1|$ in equation $y(x)$, we conclude that:

$$y(x) = c_0 \frac{1}{\sqrt{2\rho}} \frac{1}{\sqrt{\frac{1}{(1-4\alpha^2)} [e^a (\cosh 2\alpha a - 2\alpha \sinh 2\alpha a) - 1] + \frac{1}{(4\beta^2+1)} [2\beta - e^a (2\beta \sin 2\beta a + \cos 2\beta a)]}}$$

$$\left[e^{\left(\frac{1}{2}+i\sqrt{\lambda^2\rho-\frac{1}{4}}\right)x} - e^{\left(\frac{1}{2}-i\sqrt{\lambda^2\rho-\frac{1}{4}}\right)x} \right]. \tag{10}$$

Where c_o arbitrary complex number with module is one (i.e. $|c_o| = 1$).

If λ satisfies equation (8) (i.e. λ eigenvalue), then equation (10) gives eigen functions for our problem T_o (corresponding to the eigenvalue λ).

Now we determine $\max_{x \in [0,a]} |y(x)|$ and its behaviour depends on, α and β .

From

$$\left| e^{\left(\frac{1}{2}+\alpha\right)x+i\beta x} - e^{\left(\frac{1}{2}-\alpha\right)x-i\beta x} \right|^2 = 2 e^x (\cosh 2\alpha x - \cos 2\beta x),$$

We conclude that

$$\left| e^{\left(\frac{1}{2}+\alpha\right)x+i\beta x} - e^{\left(\frac{1}{2}-\alpha\right)x-i\beta x} \right| = \sqrt{2 e^x (\cosh 2\alpha x - \cos 2\beta x)}$$

Therefore,

$$|y(x)| = \left| \frac{c_o}{\sqrt{2\rho}} \frac{e^{\left(\frac{1}{2}+i\sqrt{\lambda^2\rho-\frac{1}{4}}\right)x} - e^{\left(\frac{1}{2}-i\sqrt{\lambda^2\rho-\frac{1}{4}}\right)x}}{\sqrt{\frac{1}{(1-4\alpha^2)} [e^a (\cosh 2\alpha a - 2\alpha \sinh 2\alpha a) - 1] + \frac{1}{(4\beta^2+1)} [2\beta - e^a (2\beta \sin 2\beta a + \cos 2\beta a)]}} \right|$$

Or

$$|y(x)| = \frac{1}{\sqrt{\rho}} \sqrt{\frac{e^x (\cosh 2\alpha x - \cos 2\beta x)}{\frac{1}{(1-4\alpha^2)} [e^a (\cosh 2\alpha a - 2\alpha \sinh 2\alpha a) - 1] + \frac{1}{(4\beta^2+1)} [2\beta - e^a (2\beta \sin 2\beta a + \cos 2\beta a)]}}$$

Then

$$\frac{1}{\sqrt{\rho}} \sqrt{\frac{e^x(\cosh 2ax - 1)}{\frac{1}{(1-4a^2)} [e^a(\cosh 2aa - 2a\sinh 2aa) - 1] + \frac{1}{(4\beta^2+1)} [2\beta - e^a(2\beta\sin 2\beta a + \cos 2\beta a)]}} \leq |y(x)|$$

$$\leq \frac{1}{\sqrt{\rho}} \sqrt{\frac{e^x(\cosh 2ax + 1)}{\frac{1}{(1-4a^2)} [e^a(\cosh 2aa - 2a\sinh 2aa) - 1] + \frac{1}{(4\beta^2+1)} [2\beta - e^a(2\beta\sin 2\beta a + \cos 2\beta a)]}}$$

Or

$$\sqrt{\frac{e^x(\cosh 2ax - 1)}{\rho[(1 - e^a(\cosh 2aa - 2a\sinh 2aa)) + (2\beta + e^a(2\beta + 1))]} \leq |y(x)| \leq$$

$$\sqrt{\frac{e^x(\cosh 2ax + 1)}{\rho[(e^a(\cosh 2aa - 2a\sinh 2aa) - 1) + \frac{1}{(4\beta^2+1)} (2\beta - e^a(2\beta + 1))]} .$$

Let $\max_{x \in [0, a]} |y(x)|$ be achieved at the point of x_0 , then

$$\max_{x \in [0, a]} |y(x)| = |y(x_0)|$$

$$\leq \sqrt{\frac{e^{x_0}(\cosh 2ax_0 + 1)}{\rho[(e^a(\cosh 2aa - 2a\sinh 2aa) - 1) + \frac{1}{(4\beta^2+1)} (2\beta - e^a(2\beta + 1))]}}$$

$$\leq \sqrt{\frac{e^a(\cosh 2aa + 1)}{\rho[(e^a(\cosh 2aa - 2a\sinh 2aa) - 1) + \frac{1}{(4\beta^2+1)} (2\beta - e^a(2\beta + 1))]}}$$

(Since e^x and $\cosh 2\alpha x$ are monotonic increasing on $[0, a]$), on the other hand

$$|y(x_o)| = \max_{x \in [0, a]} |y(x)| \geq |y(a)| \geq$$

$$\sqrt{\frac{e^a(\cosh 2\alpha a - 1)}{\rho[(1 - e^a(\cosh 2\alpha a - 2\alpha \sinh 2\alpha a)) + (2\beta + e^a(2\beta + 1))]}}$$

Therefore

$$\sqrt{\frac{e^a(\cosh 2\alpha a - 1)}{\rho[(1 - e^a(\cosh 2\alpha a - 2\alpha \sinh 2\alpha a)) + (2\beta + e^a(2\beta + 1))]}} \leq \max_{x \in [0, a]} |y(x)|$$

$$\leq \sqrt{\frac{e^a(\cosh 2\alpha a + 1)}{\rho \left[(e^a(\cosh 2\alpha a - 2\alpha \sinh 2\alpha a) - 1) + \frac{1}{(4\beta^2 + 1)}(2\beta - e^a(2\beta + 1)) \right]}}$$

Or

$$\sqrt{\frac{(\cosh 2\alpha a - 1)}{\rho \left[\left(\frac{1}{e^a} - (\cosh 2\alpha a - 2\alpha \sinh 2\alpha a) \right) + \left(\frac{2\beta}{e^a} + (2\beta + 1) \right) \right]}} \leq \max_{x \in [0, a]} |y(x)|$$

$$\leq \sqrt{\frac{(\cosh 2\alpha a + 1)}{\rho \left[\left((\cosh 2\alpha a - 2\alpha \sinh 2\alpha a) - \frac{1}{e^a} \right) + \frac{1}{(4\beta^2 + 1)} \left(\frac{2\beta}{e^a} - (2\beta + 1) \right) \right]}} \tag{11}$$

Now, in the obtained equation (11) used parameters α and β clearly are not parts of the equation (1) and the boundary and normalized conditions (2)-(4). Therefore we express α and β through ρ .

Suppose $\arg \lambda = \theta$, then $\lambda^2 = |\lambda|^2 (\cos 2\theta + i \sin 2\theta)$.

$$\lambda^2 \rho - \frac{1}{4} = \rho |\lambda|^2 \cos 2\theta - \frac{1}{4} + i \rho |\lambda|^2 \sin 2\theta$$

On the other hand $-\left(\lambda^2 \rho - \frac{1}{4}\right) = (\alpha + i\beta)^2 = \alpha^2 - \beta^2 + i2\alpha\beta,$

Hence

$$\alpha^2 - \beta^2 = -\rho|\lambda|^2 \cos 2\theta + \frac{1}{4}$$

$$2\alpha\beta = -\rho|\lambda|^2 \sin 2\theta$$

Or

$$\alpha^2 - \beta^2 = -\rho|\lambda|^2 \cos 2\theta + \frac{1}{4}$$

$$4\alpha^2\beta^2 = \rho^2|\lambda|^4 \sin^2 2\theta$$

Solving these two last systems of equations, we get

$$\alpha^2 = \frac{-\rho|\lambda|^2 \cos 2\theta + \frac{1}{4} + \sqrt{\left(\rho|\lambda|^2 \cos 2\theta - \frac{1}{4}\right)^2 + \rho^2|\lambda|^4 \sin^2 2\theta}}{2}$$

and

$$\beta^2 = \frac{\rho^2|\lambda|^4 \sin^2 2\theta}{2\left[-\rho|\lambda|^2 \cos 2\theta + \frac{1}{4} + \sqrt{\left(\rho|\lambda|^2 \cos 2\theta - \frac{1}{4}\right)^2 + \rho^2|\lambda|^4 \sin^2 2\theta}\right]}$$

(Since $\alpha^2 \geq 0$, then chose non negative root). Separating out the factor $\rho|\lambda|^2$ from the last relations, we deduce

$$\alpha^2 = \rho|\lambda|^2 \left(\frac{-\cos 2\theta + \frac{1}{4\rho|\lambda|^2} + \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}}{2} \right)$$

and

$$\beta^2 = \frac{\rho |\lambda|^2 \sin^2 2\theta}{2 \left[-\cos 2\theta + \frac{1}{4\rho|\lambda|^2} + \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2} \right]}$$

Or

$$\alpha = |\lambda| \sqrt{\frac{-\rho \cos 2\theta + \frac{1}{4|\lambda|^2} + \rho \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}}{2}}$$

and

$$\beta = \frac{\sqrt{\rho} |\lambda| \sin 2\theta}{\sqrt{-2\cos 2\theta + \frac{1}{2\rho|\lambda|^2} + 2 \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}}}$$

(We take the positive root and for negative root we proceed by similar way).

By substituting α and β in equation (11) and making some algebraic operations we get:

$$\frac{1}{\sqrt{|\lambda|}} \left(\frac{\cosh 2a \left(|\lambda| \sqrt{\frac{-\rho \cos 2\theta + \frac{1}{4|\lambda|^2} + \rho \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}}{2}} \right) - 1}{\frac{\rho e^{-a}}{|\lambda|} - \frac{\rho}{|\lambda|} \cosh 2a \left(|\lambda| \sqrt{\frac{-\rho \cos 2\theta + \frac{1}{4|\lambda|^2} + \rho \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}}{2}} \right)} \right.$$

$$+ \rho \sqrt{2} \sqrt{-\rho \cos 2\theta + \frac{1}{4|\lambda|^2} + \rho \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}}$$

$$\left. \sinh 2a \left(|\lambda| \sqrt{\frac{-\rho \cos 2\theta + \frac{1}{4|\lambda|^2} + \rho \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}}{2}} \right) \right.$$

$$\frac{2e^{-a} \sqrt{\rho^3 \sin 2\theta}}{\sqrt{-2 \cos 2\theta + \frac{1}{2\rho|\lambda|^2} + 2 \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}}} +$$

$$\frac{2 \sqrt{\rho^3 \sin 2\theta}}{\sqrt{-2 \cos 2\theta + \frac{1}{2\rho|\lambda|^2} + 2 \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}}} + \frac{\rho}{|\lambda|}$$

$$\leq \max_{x \in [0, a]} |y(x)| \leq$$

$$\frac{1}{\sqrt{|\lambda|}} \left(\begin{array}{l} \cosh 2a \left(|\lambda| \sqrt{\frac{-\rho \cos 2\theta + \frac{1}{4|\lambda|^2} + \rho}{2} \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}} \right) + 1 \\ \frac{\rho}{|\lambda|} \cosh 2a \left(|\lambda| \sqrt{\frac{-\rho \cos 2\theta + \frac{1}{4|\lambda|^2} + \rho}{2} \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}} \right) \\ -\rho \sqrt{2} \sqrt{-\rho \cos 2\theta + \frac{1}{4|\lambda|^2} + \rho} \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2} * \\ \sinh 2a \left(|\lambda| \sqrt{\frac{-\rho \cos 2\theta + \frac{1}{4|\lambda|^2} + \rho}{2} \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}} \right) - \frac{\rho e^{-a}}{|\lambda|} \\ + \frac{\rho}{|\lambda|} \left(\frac{-\rho \cos 2\theta + \frac{1}{4|\lambda|^2} + \rho \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}}{2\rho|\lambda|^2 \sin^2 2\theta - \cos 2\theta + \frac{1}{4\rho|\lambda|^2} + \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}} \right) \\ * \left(\frac{2e^{-a} \sqrt{\rho} \sin 2\theta}{\sqrt{-2 \cos 2\theta + \frac{1}{2\rho|\lambda|^2} + 2} \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}} - \frac{2|\lambda| \sqrt{\rho} \sin 2\theta}{\sqrt{-2 \cos 2\theta + \frac{1}{2\rho|\lambda|^2} + 2} \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}} - 1 \right) \end{array} \right)$$

Let $K_1 =$

$$\begin{aligned}
 & \cosh 2a \left(|\lambda| \sqrt{\frac{-\rho \cos 2\theta + \frac{1}{4|\lambda|^2} + \rho \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}}{2}} \right) - 1 \\
 & \frac{\rho e^{-a}}{|\lambda|} - \frac{\rho}{|\lambda|} \cosh 2a \left(|\lambda| \sqrt{\frac{-\rho \cos 2\theta + \frac{1}{4|\lambda|^2} + \rho \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}}{2}} \right) \\
 & + \rho \sqrt{2} \sqrt{-\rho \cos 2\theta + \frac{1}{4|\lambda|^2} + \rho \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}} * \\
 & \sinh 2a \left(|\lambda| \sqrt{\frac{-\rho \cos 2\theta + \frac{1}{4|\lambda|^2} + \rho \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}}{2}} \right) \\
 & \frac{2e^{-a} \sqrt{\rho^3 \sin 2\theta}}{\sqrt{-2 \cos 2\theta + \frac{1}{2\rho|\lambda|^2} + 2 \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}} + \\
 & \frac{2 \sqrt{\rho^3 \sin 2\theta}}{\sqrt{-2 \cos 2\theta + \frac{1}{2\rho|\lambda|^2} + 2 \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}} + \frac{\rho}{|\lambda|}
 \end{aligned}$$

And

$K_2 =$

$$\begin{aligned}
 & \cosh 2a \left(|\lambda| \sqrt{\frac{-\rho \cos 2\theta + \frac{1}{4|\lambda|^2} + \rho \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}}{2}} \right) + 1 \\
 & \frac{\rho}{|\lambda|} \cosh 2a \left(|\lambda| \sqrt{\frac{-\rho \cos 2\theta + \frac{1}{4|\lambda|^2} + \rho \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}}{2}} \right) \\
 & -\rho \sqrt{2} \sqrt{-\rho \cos 2\theta + \frac{1}{4|\lambda|^2} + \rho \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}} * \\
 & \sinh 2a \left(|\lambda| \sqrt{\frac{-\rho \cos 2\theta + \frac{1}{4|\lambda|^2} + \rho \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}}{2}} \right) - \frac{\rho e^{-a}}{|\lambda|} \\
 & + \frac{\rho}{|\lambda|} \left(\frac{-\rho \cos 2\theta + \frac{1}{4|\lambda|^2} + \rho \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}}{2\rho|\lambda|^2 \sin^2 2\theta - \cos 2\theta + \frac{1}{4\rho|\lambda|^2} + \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}} \right) \\
 & * \left(\frac{2e^{-a} \sqrt{\rho} \sin 2\theta}{\sqrt{\frac{-2 \cos 2\theta + \frac{1}{2\rho|\lambda|^2} + 2 \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}}{2|\lambda| \sqrt{\rho} \sin 2\theta}} - 1 \right) \\
 & \sqrt{\frac{-2 \cos 2\theta + \frac{1}{2\rho|\lambda|^2} + 2 \sqrt{1 - \frac{1}{2\rho|\lambda|^2} \cos 2\theta + \left(\frac{1}{4\rho|\lambda|^2}\right)^2}}{2|\lambda| \sqrt{\rho} \sin 2\theta}} - 1
 \end{aligned}$$

Then

$$\frac{1}{\sqrt{|\lambda|}} K_1 \leq \max_{x \in [0, a]} |y(x)| \leq \frac{1}{\sqrt{|\lambda|}} K_2.$$

Thus the proof of theorem is completed.

REFERENCES

- [1] Aigunov G.A and Jwamer K.H, Asymptotic behavior of orthonormalized eigenfunctions in a Regge type problem with a summable positive weight function, (2009), UMN , Moscow , Vol(64)6, P.169-170.
- [2] Aigounov G.A, Jwamer K. H and Dzhalaeva G.A, Estimates for the eigenfunctions of the Regge problem, Mathematical Notes ,(2012), Vol. 92, No.7, pp.127-130.
- [3] Aigunov, G. A, the boundedness of the orthonormal eigenfunctions of a certain class of non-linear Sturm-Liouville type operators with a weight function of unbounded variation on a finite interval ,Russian Math, Surveys, (2000), Vol. 55, No. 4, pp. 815-821.
- [4] Gadzhieva, T. Yu, Analysis of spectral characteristics of one non self adjoint problem with smooth coefficients, PhD thesis, Dagestan State university, (2010), South of Russian.
- [5] Jwamer K.H and Aigounov G.A., About Uniform Limitation of Normalized Eigenfunctions of T.Regge Problem in the Case of Weight Functions, Satisfying to Lipschitz Condition, Gen. Math. Notes, (2010) 1(2), 115-129.
- [6] Jwamer K. H and Qadir K.H , Estimation of Normalized eigenfunctions of spectral problem with smooth coefficients, Acta Universitatis Apulensis, Special Issue, Romania, (2011), P.113-132.
- [7] Jwamer K. H and Qadir K.H., Estimates Normalized Eigenfunction to the Boundary Value Problem in Different Cases of Weight Functions, Int. J. Open Problems Compt.Math.,(2011), Vol. 4(3), P.62-71 .
- [8] Kravitsky A.O , On series expansion in eigen functions of one non self-adjoint boundary problem, Report of Academy of Science, USSR, (1966), Vol.170, No.6, P.1255-1258.
- [9] Naimark. M. A, Linear differential operators, 2nd edition, Nauka, Moscow, (1969), English trans 1. Of 1st edition, Vols. I, II, Ungar, New York, 1967, 1968.
- [10] Regge .T , Analytical properties of the scattering matrix , Mathematics(collection of translations), (1963), Vol.4, P.83-89.